

Quantum corrected full-band Cellular Monte Carlo simulation of AlGa_N/Ga_N HEMTs

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A full-band Cellular Monte Carlo (CMC) approach [1] is applied to simulation of electron transport in AlGa_N/Ga_N HEMTs with quantum corrections included via the effective potential method. The full-band CMC transport model is based on a detailed model of the electron-phonon interactions in the wurtzite crystal structure using the rigid pseudo-ion model, where the anisotropic deformation potentials are derived from the electronic band structure, the atomic pseudopotential, and the phonon dispersion. Realistic polar-optical phonon, impurity, piezoelectric and dislocation scatterings are also included in the full-band CMC simulator, which shows good agreement with measured velocity-field data from pulsed I-V measurements at Arizona State University [2].

AlGa_N/Ga_N HEMTs have several advantages over GaAs based technology due to their potential for high-voltage and high-power operation. An important issue of AlGa_N/Ga_N HEMTs is the high sheet charge density N_s , which is obtained even without modulation doping in the AlGa_N layer due to the piezoelectric and spontaneous polarization. In order to include quantum effects into the transport simulation, one needs to solve the two-dimensional Schrödinger-Poisson problem with an appropriate transport kernel. An alternative way is the use of effective potentials [3] with a particle-based Monte Carlo simulator. The effective potential, $V_{eff}(x)$, is obtained by the integral transformation from the classical potential $V(x)$ as

$$V_{eff}(x) = \frac{1}{\sqrt{2\pi}a_0} \int_{-\infty}^{+\infty} V(x + \xi) \exp\left(-\frac{\xi^2}{2a_0^2}\right) d\xi . \quad (1)$$

The calculated HEMT structure consists of a 15 nm doped Al_{0.2}Ga_{0.8}N layer, 5 nm unintentionally doped Al_{0.2}Ga_{0.8}N layer, and 100 nm unintentionally doped Ga_N layer, shown in Fig. 2. The background, unintentional doping is taken as 10^{17} cm⁻³ for both the Ga_N and AlGa_N layers. The electron densities with various Gaussian smoothing parameters, a_0 , are shown in Fig. 3. The Schrödinger-Poisson result is also shown in the figure. This figure shows that the use of $a_0=3\text{Å}$ can quite accurately describe the electron density reduction and charge setback from the interface.

This effective potential approach has been combined with the full-band CMC code for device simulation. Figure 4 shows the electron distribution along the channel. The electron displacement due to the effective potential can be clearly seen in the figure. The Calculated I_d - V_{ds} characteristics, with and without the effective potential, are shown in Fig. 5. An approximately 20 % reduction due to quantum-mechanical size-quantization effects is found in the current.

[1] M. Saraniti and S.M. Goodnick, *IEEE Trans. Elec. Dev.* **47**, 1909 (2000)

[2] J.M. Barker *et al.*, *Physica* **B314**, 39 (2002)

[3] D.K. Ferry, *Superlattices and Microstructures* **28**, 419 (2000)

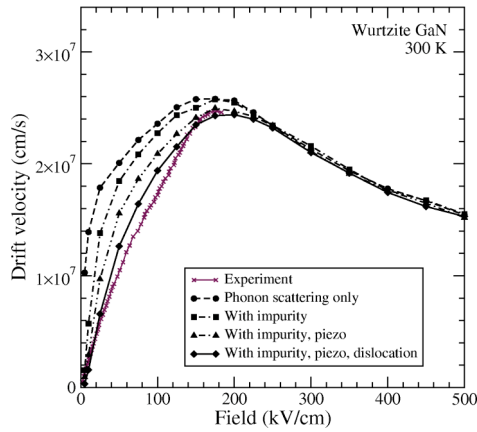


Fig. 1 : Verocity vs field relationship for bulk GaN.

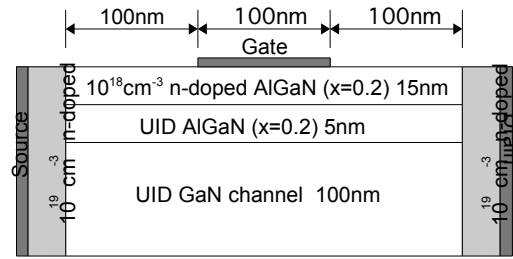


Fig. 2 : Structure of the calculated HEMT device.

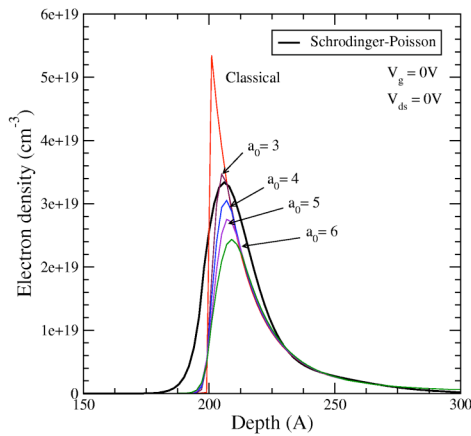


Fig. 3 : Electron density by the classical and effective potential method. The electron density by the Schrödinger-Poisson calculation is also shown.

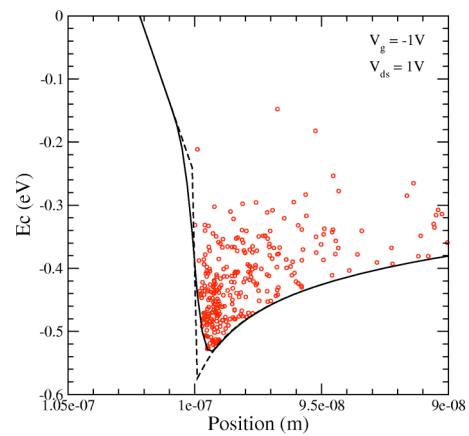


Fig. 4 : Electron distribution along the channel using the effective potential with $a_0=3\text{Å}$. Note the displacement of electrons by the effective potential.

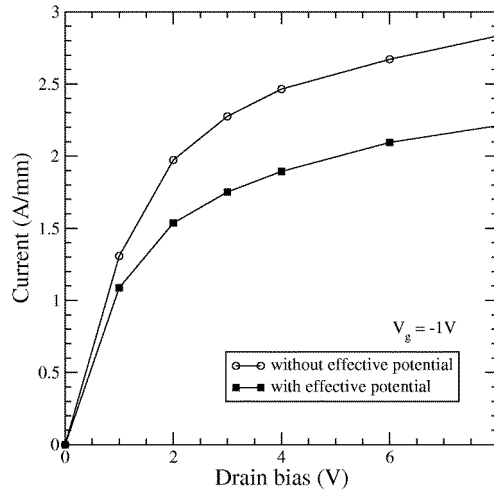


Fig. 5 : Output device characteristics with and without the effective potential. The gate voltage is $-1V$. Notice that there is about 20 % current reduction due to the quantization effect.